

MRC Technical Summary Report #1875

REMARKS ON THE SCHRÖDINGER OPERATOR
WITH SINGULAR COMPLEX POTENTIALS

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August 1978

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(Received June 7, 1978)



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UNIVERSITY OF WISCONSIN - MADISON MATHEMATICS RESEARCH CENTER REMARKS ON THE SCHRÖDINGER OPERATOR WITH SINGULAR COMPLEX POTENTIALS Haim Brezis 9 9 Tosio Kato 94 Technical summary kep 1875 (12) 22 P ABSTRACT

Let $A = -\Delta + V(x)$ be a Schrödinger operator on an (arbitrary) open set $\Omega \subseteq \mathbb{R}^m$, where $V \in L^1_{loc}(\Omega)$ is a complex valued function. We consider the "maximal" realization of A in $L^2(\Omega)$ under Dirichlet boundary condition, that is

$$D(A) = \{u \in H_0^1(\Omega); \ Vu \in L_{loc}^1(\Omega) \ \text{and} \ -\Delta u + Vu \in L^2(\Omega)\} \ .$$

When $\Omega = \mathbb{R}^m$ we also consider the operator (14) MR2-T5R-1875

 $A_{\star} = -\Delta + V$

with domain

$$D(A_1) = \{ u \in L^2(\Omega); \ Vu \in L^1_{loc}(\Omega) \ \text{and} \ -\Delta u + Vu \in L^2(\Omega) \} \ .$$

A special case of our main results is:

Theorem: Let $m \ge 3$; assume that the function $max\{-Re\ V,0\}$ belongs to $L^{\infty}(\Omega) + L^{m/2}(\Omega)$ and also to $L^{(m/2)+\epsilon}_{loc}(\Omega)$ for some $\epsilon > 0$. Then A (resp. A₁) is closable and $\tilde{A} + \lambda$ (resp. $\tilde{A}_1 + \lambda$) is m-accretive for some real constant λ .

AMS (MOS) Subject Classifications: 35J10, 47B44

Key Words: Schrödinger operator, Complex potentials, m-accretive operator

Work Unit Number 1 (Applied Analysis)

DAAG29-75-C-00245NSF-MC576-046551

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- 3) the United States Army under Contract No. DAAG29-75-C-0024;
- 4) the National Science Foundation under Grant No. MCS76-04655.

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SIGNIFICANCE AND EXPLANATION

Schrödinger operators of the form A = -0 + V(x), where 0 is the Laplacian and V is a scalar potential, arise in quantum mechanics and other areas. Delicate questions concerning what domain should be assigned to A must be settled in order to have a good theory. These questions are answered here for a very general class of potentials V which may even have complex values.

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REMARKS ON THE SCHRÖDINGER OPERATOR WITH SINGULAR COMPLEX POTENTIALS Haim Brezis1,3 and Tosio Kato2,4

1. Introduction

Let $A = -\Delta + V(x)$ be a Schrödinger operator on an (arbitrary) open set $\Omega \subset \mathbb{R}^m$, where $V \in L^1_{loc}(\Omega)$ is a complex valued function. We consider the "maximal" realization of A in $L^2(\Omega)$ under Dirichlet boundary condition, that is

$$D(A) = \{ u \in H_0^1(\Omega); \ Vu \in L_{loc}^1(\Omega) \ \text{ and } \ -\Delta u + Vu \in L^2(\Omega) \} \ .$$

When $\Omega = \mathbb{R}^m$ we also consider the operator

$$A_1 = -\Delta + V$$

with domain

$$D(A_1) = \{ u \in L^2(\Omega) ; \ Vu \in L^1_{loc}(\Omega) \ \text{ and } -\Delta u + Vu \in L^2(\Omega) \} \ .$$

We state now our main results (see Theorems 3.1 and 3.2) in a special case.

Theorem: Let $m \ge 3$; assume that the function $\max\{-\text{Re V}, 0\}$ belongs to $L^{\infty}(\Omega) + L^{m/2}(\Omega)$ and also to $L_{\text{loc}}^{(m/2)+\epsilon}(\Omega)$ for some $\epsilon>0$. Then A (resp. A_1) is closable and $\bar{A}+\lambda$ (resp. $\tilde{A}_1 + \lambda$) is m-accretive for some real constant λ .

We emphasize the fact that max(Re V, 0) and ImV could be arbitrary functions in $L_{loc}^{1}(\Omega)$.

Our methods rely on some measure theoretic arguments and standard techniques of DeGiorgi-Moser-Stampacchia type, related to the weak form of the maximum principle.

The distributional inequality

proved in [3] plays a crucial role. We also use a result from [1] concerning a property of Sobolev spaces.

In order to describe our method in a simple case we begin in Section 2 with real valued potentials. The main results in Section 2 are essentially known (see [3], [4], [8]) except perhaps for Theorem 2.2 when m < 4.

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In Section 3 we turn to the case of complex potentials. Schrödinger operators with complex potentials have been studied by Nelson [6]. His results were extended in [5]. Here we allow more general singularities.

We thank Professors R. Jensen and B. Simon for useful suggestions and discussions (with the first author).

2. Real valued potentials

Let Ω be an (arbitrary) open subset of \mathbb{R}^m and let $H=L^2=L^2(\Omega;\mathfrak{C})$. Let $q\in L^1_{loc}(\Omega)$ be a real valued function. Set

$$q^+ = \max(q,0), \quad q^- = \max(-q,0)$$
.

Assume

(1) $q \in L^{\infty}(\Omega) + L^{p}(\Omega)$

with

$$\begin{cases} p = \frac{m}{2} & \text{when} & m \ge 3 \\ p > 1 & \text{when} & m = 2 \\ p = 1 & \text{when} & m = 1 \end{cases}$$

Consider the operator A defined in H by

$$A = -\Delta + q(x)$$

with

$$D(A) = \{u \in H_0^1(\Omega); qu \in L_{loc}^1(\Omega) \text{ and } -\Delta u + qu \in L^2(\Omega)\}.$$

The main results are the following:

Theorem 2.1. A is self-adjoint and $\lambda + \lambda_1$ is m-accretive for some real constant λ_1 . Furthermore $u, v \in D(\lambda)$ imply $q|u|^2 \in L^1(\Omega)$, $q|v|^2 \in L^1(\Omega)$ and

When $\Omega = \mathbb{R}^{\mathbb{R}}$ we also consider the operator λ_1 defined in H by

$$A_1 = -\Delta + q(x)$$

with

$$D(A_1) = \{u \in L^2(\Omega); qu \in L^1_{loc}(\Omega) \text{ and } -\Delta u + qu \in L^2(\Omega)\}.$$

Only when m = 3 or m = 4 we will make the additional assumption:

(3) $q^{-\epsilon} L_{loc}^{p+\epsilon}(\Omega)$ with $p=\frac{3}{2}$ when m=3 and p=2 when m=4, for some arbitrarily small $\epsilon>0$.

More precisely we assume that for each $x_0 \in \mathbb{R}^m$ there exists a neighborhood U of x_0 and some $\epsilon > 0$ (depending on x_0) such that $q^- \in L^{p+\epsilon}(U)$.

Theorem 2.2: Under the assumptions (1) and (3), $\lambda_1 = \lambda$.

Our first lemma is well known:

Lemma 2.1: Assume (1). Then for every $\varepsilon > 0$, there exists a constant λ_{ε} such that

$$\int q^{-} \left| u \right|^{2} \leq \varepsilon \left\| \operatorname{gradu} \right\|_{L^{2}}^{2} + \lambda_{\varepsilon} \left\| u \right\|_{L^{2}}^{2} \qquad \forall u \in \operatorname{H}_{0}^{1}(\Omega) \ .$$

In particular

$$\int q |\mathbf{u}|^2 \leq \left\| \operatorname{gradu} \right\|_{\mathbf{L}^2}^2 + \lambda_1 \left\| \mathbf{u} \right\|_{\mathbf{L}^2}^2 \qquad \forall \mathbf{u} \in \mathbf{H}_0^1(\Omega) \ .$$

<u>Proof:</u> Write $q = q_1 + q_2$ with $q_1 \in L^{\infty}(\Omega)$ and $q_2 \in L^{p}(\Omega)$. Then for each k > 0 we have

$$\int q^{-}|u|^{2} \leq ||q_{1}||_{L^{\infty}} ||u||_{L^{2}}^{2} + \int_{\{|q_{2}|>k\}} |q_{2}||u|^{2} + k \int_{\{|q_{2}|\leq k\}} |u|^{2}$$

$$\leq (\|q_1\|_{L^{\infty}} + k) \|u\|_{L^2}^2 + \|q_2\|_{L^{p}([q_2] > k])} \|u\|_{L^2}^2$$

with

$$\frac{1}{p} + \frac{2}{t} = 1.$$

In case $m \ge 3$ we find $t = 2^{\frac{n}{2}}$ where $2^{\frac{n}{2}}$ is the Sobolev exponent, that is $\frac{1}{2^{\frac{n}{2}}} = \frac{1}{2} - \frac{1}{m}$. By the Sobolev imbedding theorem we have

$$\|\mathbf{u}\|_{\mathbf{L}^{\frac{1}{2}}} \leq c \|\mathbf{gradu}\|_{\mathbf{L}^{2}} \qquad \forall \mathbf{u} \in \mathbf{H}_{0}^{1}(\Omega) \ .$$

When m = 2 we find $2 < t < \infty$ and it is known that

$$\left\| \mathbf{u} \right\|_{\mathbf{L}^{\xi}} \leq C \left(\left\| \mathbf{gradu} \right\|_{\mathbf{L}^{2}} + \left\| \mathbf{u} \right\|_{\mathbf{L}^{2}} \right) \qquad \forall \mathbf{u} \in \mathbf{H}_{0}^{1}(\Omega) \ .$$

When m = 1 we find $t = \infty$ and it is known that

$$\left\| \mathbf{u} \right\|_{\mathbf{L}^{\infty}} \leq C(\left\| \operatorname{gradu} \right\|_{\mathbf{L}^{2}} + \left\| \mathbf{u} \right\|_{\mathbf{L}^{2}}) \qquad \forall \mathbf{u} \in \mathbf{H}_{0}^{1}(\Omega) \ .$$

We reach the conclusion of Lemma 2.1 in all the cases by choosing k large enough so that

$$c^{2}\|q_{2}\|_{L^{p}([|q_{2}|>k])} < \varepsilon$$
.

Remark 2.1: Assumption (1) is used in all the results of this paper only through

Lemma 2.1 and it may in fact be weakened to a "locally uniform LP-condition":

(1')
$$||q^{-}||_{L^{\mathbf{p}}(\Omega \cap B_{\mu}(y))} \to 0 \text{ as } r \to 0 \text{ uniformly in } y \in \Omega ,$$

where

$$B_r(y) = \{x \in \mathbb{R}^m; |x - y| \le r\}$$
.

Indeed let $\varphi \in \mathcal{D}_+(\mathbb{R}^m)$ with supp $\varphi \subseteq \mathbb{B}_{\mathbb{R}}(0)$ and $||\varphi||_{\mathbb{L}^2} = 1$. Then, writing $\varphi_{\mathbb{V}}(\mathbf{x}) = \varphi(\mathbf{x} - \mathbf{y})$,

$$\int q^{-}|u|^{2} = \int dy \int q^{-}|u\varphi_{y}|^{2} \le \int ||q^{-}||_{L^{p}(B_{x}(y))} ||u\varphi_{y}||_{L^{t}}^{2} dy .$$

Here $\|q^-\|_{L^p(B_r(y))} \le \delta$ for any small δ by (1') if r is chosen small. So

Choosing δ so that $2C\delta = \epsilon$, one gets the conclusion of Lemma 2.1. Such a locally uniform L^p -condition was used by Simader [7].

We recall a result of [1] which will be used in the proof of Theorem 2.1 (1).

Lemma 2.2: Let $T \in H^{-1}(\Omega) \cap L^1_{loc}(\Omega)$ and let $u \in H^1_0(\Omega)$ be such that a.e. on Ω

for some real valued function f ϵ L¹(Ω). Then Re T \cdot \bar{u} ϵ L¹(Ω) and Re(T,u) = \int ReT \cdot \bar{u}

where (T,u) denotes the Hermitian scalar product in the duality between $H^{-1}(\Omega)$ and $H^1_{\Omega}(\Omega)$.

The proof of Theorem 2.1 is divided into 4 steps.

Step 1: $\lambda + \lambda$ is onto for $\lambda > \lambda_1$. Set $q_n^+ - \min(q^+, n)$; by a Theorem of Lax-Milgram there exists a unique function $u_n \in H^1_0(\Omega)$ which satisfies

(4)
$$-\Delta u_n + (q_n^+ - q_n^-)u_n + \lambda u_n = f.$$

⁽¹⁾ The use of this sort of lemma in this context was suggested by M. Crandall.

(Note that by Lemma 2.1 the sesquilinear form $\int q u v$ is continuous on $H_0^1(\Omega)$). Multiplying (4) by u we find a constant C independent of n such that

(5)
$$\|u_n\|_{H^{1}} \leq c,$$

$$\int q_n^+ |u_n|^2 \le c.$$

Choose a subsequence denoted again by u_n such that $u_n + u$ weakly in $H_0^1(\Omega)$ and $u_n + u$ a.e. on Ω . It follows from Faton's Lemma and (6) that $q^+|u|^2 \epsilon L^1(\Omega)$. We deduce that $qu \in L^1_{loc}(\Omega)$; indeed

$$q^{+}|u| \leq \frac{1}{2} q^{+}(|u|^{2} + 1) \in L^{1}_{loc}(\Omega) ,$$

 $q^{-}|u| \leq \frac{1}{2} q^{-}(|u|^{2} + 1) \in L^{1}_{loc}(\Omega) .$

We pass now to the limit in (4) and prove that $-\Delta u + qu + \lambda u = f$ in $D'(\Omega)$. It suffices to show that

$$(q_n^+ - q^-)u_n + qu$$
 in $L^1_{loc}(\Omega)$.

For this purpose we adapt a device due to W. Strauss [9] and extensively used in the study of strongly nonlinear equations. In view of Vitali's convergence theorem, it suffices to verify that given $\omega \subset \Omega$, then $\forall \varepsilon > 0$, $\exists \delta > 0$ such that $E \subset \omega$ and $|E| < \delta$ imply $\int_E |q_n^+ - q^-||u_n^-|| < \varepsilon$ for all n. But for every R > 0 we have

$$q_n^+|u_n^-| \le \frac{1}{2} q_n^+(R + \frac{1}{R} |u_n^-|^2)$$

and thus, by (6),

$$\int_{E} q_{n}^{+} |u_{n}^{-}| \leq \frac{1}{2} R \int_{E} q^{+} + \frac{1}{2R} C .$$

We fix R large enough so that $\frac{C}{R} < \epsilon$ and then $\delta > 0$ so small that R $\int_{E} q^{+} < \epsilon$. We proceed similarly with $q^{-}|u_{n}|$.

Step 2: $A + \lambda_1$ is accretive. Let $u \in D(A)$ and set T = qu. Since $T \in H^{-1}(\Omega) \cap L^1_{loc}(\Omega)$ and

Re
$$Tu = q|u|^2 \ge -q^-|u|^2 \in L^1(\Omega)$$

it follows from Lemma 2.2 that $q|u|^2 \in L^1$ and

$$Re(T,u) = \int q|u|^2$$
.

But qu = Au + Au and so

$$Re(Au,u) - \int |gradu|^2 = \int g|u|^2$$
.

Since Au $\in L^2(\Omega)$ we have in fact

Re(Au,u) =
$$\int |gradu|^2 + \int q|u|^2 \ge -\lambda$$
, $\int |u|^2$

by Lemma 2.1.

Step 3: $u \in D(A)$ implies $q|u|^2 \in L^1(\Omega)$ and (2) holds. We have just seen in Step 2 that $u \in D(A)$ implies $q|u|^2 \in L^1(\Omega)$. Now let $u, v \in D(A)$ and set T = qu. We have $T \in H^{-1}(\Omega) \cap L^1_{loc}(\Omega)$ and

Re T ·
$$\vec{v}$$
 = Re $qu\vec{v} \ge -\frac{1}{2}|q||u|^2 - \frac{1}{2}|q||v|^2 \in L^1(\Omega)$

and therefore

$$Re(T,v) = \int Re quv$$
.

Thus

Changing u into iu we find

$$(Au,v) = \int gradu gradv + \int quv$$
.

<u>Step 4</u>: A is self-adjoint. Indeed A + λ_1 is m-accretive and symmetric. Therefore A + λ_1 is self-adjoint and so is A.

<u>Proof of Theorem 2.2</u>: Clearly $A \subseteq A_1$. Let $u \in D(A_1)$ and set $f = A_1u + \lambda u$ with some $\lambda > \lambda_1$. Let $u \in D(A)$ be the unique solution of

$$Au'' + \lambda u'' = f.$$

We have

$$A_1(u - u^*) + \lambda(u - u^*) = 0$$
.

Since $(u - u^*) \in L^1_{loc}(\mathbb{R}^m)$ and $\Delta(u - u^*) \in L^1_{loc}(\mathbb{R}^m)$ we may apply Lemma A in [3] to conclude that

$$\Delta |u - u^*| \ge \text{Re}[\Delta(u - u^*) \text{ sign}(\bar{u} - \bar{u}^*)] \text{ in } \mathcal{D}^*(\mathbb{R}^m)$$
,

and thus in $D'(\mathbb{R}^{\mathbb{R}})$ we find,

$$\Delta |u - u^*| \ge Re[(q + \lambda)|u - u^*|] \ge (-q^- + \lambda)|u - u^*|$$
.

Using the next lemma we conclude that $u = u^*$ (and hence $D(A_1) = D(A)$).

<u>Lemma 2.3</u>: Assume (1) and (3). Let $v \in L^2(\mathbb{R}^m)$ be a real valued function with $q^{-}v \in L^1_{loc}(\mathbb{R}^m)$ satisfying

$$-\Delta v - q v + \lambda v \le 0$$
 in $D'(\mathbb{R}^m)$

with some $\lambda > \lambda_1$. Then $v \le 0$ a.e. on \mathbb{R}^m .

The proof of Lemma 2.3 relies on the following crucial result. Since we shall need it in Section 3 for a general domain $\Omega \subset \mathbb{R}^m$ we work now again in Ω .

Theorem 2.3: Assume (1). Let $g \in L^2(\Omega) \cap L^\infty(\Omega)$ and let $\psi \in H^1_0(\Omega)$ be the unique solution of

(7)
$$-\Delta \psi - q \psi + \lambda \psi = g \text{ in } \Omega \quad (\lambda > \lambda_1) .$$

Then

a) $g \ge 0$ a.e. on Ω implies $\psi \ge 0$ a.e. on Ω ;

b)
$$\psi \in \bigcap_{2 \le p < \infty} L^p(\Omega)$$
.

Proof of Theorem 2.3: a) Multiplying (7) by -\u03c4 we find

$$\int |grad \psi^{-}|^{2} - \int q^{-} |\psi^{-}|^{2} + \lambda \int |\psi^{-}|^{2} \leq 0$$

and thus $\psi = 0$.

b) We have to consider only the case $m \ge 3$ (when $m \le 2$, $\psi \in H^1_0(\Omega)$ implies $\psi \in \bigcap_{2 \le p \le m} L^p(\Omega)$).

We can always assume that $g \ge 0$ a.e. on Ω so that $\psi \ge 0$ a.e. on Ω . We truncate q by $q_k = \min(q^-, k)$ and define ψ_k to be the unique solution of

$$\begin{cases} \psi_{\mathbf{k}} \in H_{0}^{1}(\Omega) \\ -\Delta \psi_{\mathbf{k}} - q_{\mathbf{k}}^{-} \psi_{\mathbf{k}} + \lambda \psi_{\mathbf{k}} = g \quad \text{in } \Omega . \end{cases}$$

It is clear that $\psi_k + \psi$ weakly in $H_0^1(\Omega)$ as $k + \infty$. We shall prove that for every $p \in (2,\infty)$, $\psi_k \in L^p(\Omega)$ and

(8)
$$\|\psi_{\mathbf{k}}\|_{\mathbf{L}^{\mathbf{p}}} \leq c_{\mathbf{p}}(\|\mathbf{g}\|_{\mathbf{L}^{2}} + \|\mathbf{g}\|_{\mathbf{L}^{\infty}})$$
,

where C_p is independent of k, but it depends on q through the use of Lemma 2.1. For simplicity we drop now the subscript k on ψ_k and write

$$(9) -\Delta \psi - q_{\nu}^{-}\psi + \lambda \psi = g .$$

Set $\psi_n = \min(\psi, n)$ and let $2 \le p < \infty$; since $(\psi_n)^{p-1} \in H_0^1(\Omega)$ we can multiply (9) by $(\psi_n)^{p-1}$ and we get

$$(p-1) \int (\psi_n)^{p-2} |\operatorname{grad} \psi_n|^2 \le \int g(\psi_n)^{p-1} + \int q_k^-(\psi_n)^p + \int |\psi_n|^{p-1} \psi ,$$

that is

$$\begin{split} &\frac{4(p-1)}{p^2} \int |\text{grad } \psi_n^{p/2}|^2 \leq \|g\|_{L^p} \|\psi_n\|_{L^p}^{p-1} + \int q^-(\psi_n)^p + kn^{p-1} \int_{\{\psi^> n\}} \psi \\ &\leq \|g\|_{L^p} \|\psi_n\|_{L^p}^{-p-1} + \varepsilon \|\text{grad} \psi_n^{p/2}\|_{L^2}^2 + \lambda_\varepsilon \|\psi_n\|_{L^p}^p + k \int_{\{\psi^> n\}} \psi^p \end{split}$$

by Lemma 2.1 (here $\int\limits_{\{\psi>n\}}\psi^p$ is possibly infinite). Choosing $\epsilon>0$ small enough (for example $\epsilon=\frac{2(p-1)}{p^2}$) we see that

$$\int \left| \operatorname{grad} \, \psi_n^{p/2} \right|^2 \leq C_p [\left\| g \right\|_{L^p}^p + \left\| \psi \right\|_{L^p}^p + k \int\limits_{\left[\psi > n \right]} \psi^p]$$

where C_{p} is independent of k and n. Using Sobolev's inequality we find

(10)
$$\|\psi\|_{L^{p_2^*/2}}^p \leq c_p \left[\|g\|_{L^p}^p + \|\psi\|_{L^p}^p + k \int_{[\psi > n]} \psi^p \right].$$

<u>Assuming</u> now that $\psi \in L^p(\Omega)$ and passing to the limit in (10) as $n \to \infty$ we obtain that $\psi \in L^{p^2/2}(\Omega)$ and

$$\|\psi\|_{L^{p_2^*/2}} \le c_p[\|g\|_{L^p} + \|\psi\|_{L^p}]$$
.

Iterating this process from p = 2 we obtain finally for every $p \in [2,\infty)$

$$\|\psi\|_{L^p} \le c_p [\|g\|_{L^2} + \|g\|_{L^\infty}]$$
.

More precisely we have proved (8). The conclusion of Theorem 2.3 follows since $\psi_k \to \psi$ weakly in $H^1_0(\Omega)$ as $k \to \infty$.

Proof of Lemma 2.3: By assumption
$$q v \in L^1_{loc}(\mathbb{R}^m)$$
 and
$$\int v(-\Delta \varphi - q \varphi + \lambda \varphi) \leq 0 \quad \forall \varphi \in D_{\perp}(\mathbb{R}^m)$$

An easy density argument (smoothing by convolution) shows that

(11)
$$\int \mathbf{v}(-\Delta \varphi - \mathbf{q} \varphi + \lambda \varphi) \leq 0 \quad \forall \varphi \in H^2(\mathbb{R}^m) \cap L^{\infty}(\mathbb{R}^m), \text{ supp } \varphi \text{ compact, } \varphi \geq 0 \text{ a.e. }.$$

Fix $g \in D_+(\mathbb{R}^m)$ and let $\psi_k \in H^1(\mathbb{R}^m)$ be the unique solution of

(12)
$$-\Delta\psi_{k} - q_{k}^{\dagger}\psi_{k} + \lambda\psi_{k} = g \text{ in } \mathbb{R}^{m}.$$

We know by Theorem 2.3 that $\psi_k \ge 0$ a.e.

$$\psi_k \in \bigcap_{2 \le p < \infty} L^p(\mathbb{R}^m)$$
 with $\|\psi_k\|_{L^p} \le c_p$,

and also $\|\operatorname{grad} \psi_{\mathbf{k}}\|_{2} \leq C$. In addition we derive from (12) that

$$\psi_{\mathbf{k}} \in H^2(\mathbf{R}^{\mathbf{m}}) \cap L_{\mathrm{loc}}^{\infty}(\mathbf{R}^{\mathbf{m}})$$
.

Fix $\zeta \in D_+(\mathbb{R}^m)$ satisfying $\zeta(x) = 1$ for $|x| \le 1$ and set $\zeta_n(x) = \zeta(\frac{x}{n})$. In (11) we choose $\varphi = \psi_k \zeta_n$. Note that by (12)

$$-\Delta \varphi - q \varphi + \lambda \varphi = \zeta_n g - (\Delta \zeta_n) \psi_k - 2 \operatorname{grad} \zeta_n \operatorname{grad} \psi_k - \zeta_n \psi_k (q - q_k),$$

and therefore

$$\int \, v \zeta_n g \, \leq \frac{C}{n^2} + \frac{C}{n} + \int \, v \zeta_n \psi_k (q^- - q_k^-) \ . \label{eq:constraint}$$

First we fix n and let $k \to \infty$. We distinguish two cases:

- a) $m \geq 5$,
- b) m < 5.
- a) When $m \ge 5$ we have $q q_k + 0$ in $L_{loc}^{m/2}(\mathbb{R}^m)$. Let $p \in [2,\infty)$ be such that $\frac{1}{2} + \frac{2}{m} + \frac{1}{p} = 1$; we have

$$|\int |v\zeta_n\psi_k(q^--q_k^-)| \leq ||v||_{L^2} ||\psi_k||_{L^p} ||\zeta_n(q^--q_k^-)||_{L^{m/2}} \to 0.$$

Consequently

$$\int v\zeta_n g \leq \frac{C}{2} + \frac{C}{n} .$$

b) When m < 5 we use the assumption (3) (or (1)): $q^{-\epsilon} L_{loc}^{m/2+\epsilon}(\mathbb{R}^m)$ with some $\epsilon > 0$. It follows from (12) that ψ_k remains bounded in $W_{loc}^{2,q}(\mathbb{R}^m)$ for some $q > \frac{m}{2}$ (when $m \ge 2$) as $k \to \infty$. We conclude that ψ_k remains bounded in $L_{loc}^{\infty}(\mathbb{R}^m)$ as $k \to \infty$ (in case m = 1, ψ_k is bounded in $L_{loc}^{\infty}(\mathbb{R}^m)$. Therefore

$$\int v \zeta_n \psi_k (q - q_k) \to 0 \text{ as } k \to \infty$$

since $\|\zeta_n v(q^7 - q_k^7)\|_{q^1} \to 0$ by the dominated convergence theorem (recall that

 $q^{T}v \in L^{1}_{loc}(\mathbb{R}^{m}))$. In both cases we find

$$\int v\zeta_n^{}g \leq \frac{C}{n^2} + \frac{C}{n} \quad \forall_n .$$

As n + w we see that

$$\int vg \leq 0 \quad \forall g \in D_+(\mathbb{R}^m)$$

and therefore $v \le 0$ a.e. on \mathbb{R}^m .

Remark 2.2: The conclusion of Lemma 2.3 fails in \mathbb{R}^3 and in \mathbb{R}^4 if we do not assume (3). Ancona (personal communication) has constructed in \mathbb{R}^3 and in \mathbb{R}^4 functions $q^- \in L^{m/2}(\mathbb{R}^m)$ and $u \in L^{m/m-2}(\mathbb{R}^m) \cap L^2(\mathbb{R}^m)$ such that

$$-\Delta u - q u + u = 0 \quad \text{in} \quad D'$$

with $\|q^T\|_{L^{m/2}}$ as small as we please and $u \neq 0$.

3. Complex potentials

Let Ω be an (arbitrary) open subset of \mathbb{R}^m . Assume q(x) and q'(x) are real valued functions such that $q,q'\in L^1_{loc}(\Omega)$ and set

$$V(x) = q(x) + iq'(x).$$

We assume

(13) either
$$q' \in L^{1+\epsilon}_{loc}(\Omega)$$
 or $q' \in L^{(m/2)+\epsilon}_{loc}(\Omega)$ when $m \ge 2$

for some arbitrarily small $\varepsilon > 0$. Define

$$A = -\Delta + V(x)$$

with

$$D(A) = \{u \in H_0^1(\Omega); \ Vu \in L_{loc}^1(\Omega) \ \text{and} \ -\Delta u + Vu \in L^2(\Omega)\} \ .$$

The main results are the following

Theorem 3.1: Assume (1) and (13). Then A is closable in $L^2(\Omega)$ and $\tilde{A} + \lambda_1$ is m-accretive. In addition $u \in D(\tilde{A})$ implies that $u \in H^1_0(\Omega)$, $q|u|^2 \in L^1(\Omega)$ and

(14)
$$\operatorname{Re}(\widetilde{A}u,u) = \int \left|\operatorname{grad}u\right|^2 + \int q|u|^2.$$

Remark 3.1: In case we assume

(15)
$$|q'(x)| \leq Mq^+(x) + h(x)$$
 for a.e. $x \in \Omega$

with $h \in L^{2m/(m+2)}_{loc}(\Omega)$ and $m \ge 3$ then A is closed in $L^2(\Omega)$. (Note that (15) corresponds essentially with the assumption made in [5]). Indeed let $u_n \in D(A)$ be such that $u_n + u$ in $L^2(\Omega)$ and $Au_n + f$ in $L^2(\Omega)$. It follows from Lemma 2.1 and (14) that $u_n + u$ in $H^1_0(\Omega)$ and $\sqrt{q} u_n + \sqrt{q} u$ in $L^2(\Omega)$. From (15) we deduce easily that $Vu \in L^1_{loc}(\Omega)$ and that $-\Delta u + Vu = f$ in $D^1(\Omega)$. Therefore $u \in D(A)$ and Au = f.

When $\Omega = \mathbb{R}^m$ we consider also the operator A_1 defined in $L^2(\mathbb{R}^m)$ by

$$A_1 = -\Delta + V(x)$$

with

$$D(A_1) = \{u \in L^2(\mathbb{R}^m); \ \forall u \in L^1_{loc}(\mathbb{R}^m) \ \text{and} \ -\Delta u + \forall u \in L^2(\mathbb{R}^m)\} \ .$$

Theorem 3.2: Assume (1), (3) and (13). Then A_1 is closable and $\overline{A_1} = \overline{A}$.

In the proof of Theorem 3.1 we shall use the following

<u>Lemma 3.1</u>: Let $\mathbf{v} \in H_0^1(\Omega)$ be a real valued function. Assume (1) and $-\Delta \mathbf{v} - \mathbf{q} \cdot \mathbf{v} + \lambda \mathbf{v} \leq 0 \quad \text{in} \quad D^*(\Omega)$

with $\lambda > \lambda_1$. Then $v \le 0$ a.e. on Ω .

Proof of Lemma 3.1: We have, for every $\varphi \in D_{\perp}(\Omega)$

$$\int gradv \ gradv - \int q w + \lambda \int w \le 0$$
.

Now we use the fact (pointed out by G. Stampacchia) that $D_+(\Omega)$ is dense in $\{u \in H_0^1(\Omega); u \geq 0 \text{ a.e. on } \Omega\}$ for the H^1 norm (1) to derive that

 $\int \operatorname{grad} v \ \operatorname{grad} \psi \ - \int q^{-} v \psi \ + \ \lambda \ \int v \psi \ \leq \ 0 \quad \forall \psi \ \varepsilon \ \operatorname{H}^{1}_{0}(\Omega) \ , \quad \phi \ \geq \ 0 \ .$

Choosing $\varphi = v^+$ we obtain

$$\int |gradv^{+}|^{2} - \int g^{-}|v^{+}|^{2} + \lambda \int |v^{+}|^{2} < 0$$

and therefore $v^{\dagger} = 0$.

The proof of Theorem 3.1 is divided into five steps.

Step 1: $R(A + \lambda) \supset L^{2}(\Omega) \cap L^{\infty}(\Omega)$ for $\lambda > \lambda_{1}$.

Indeed let $f \in L^2(\Omega) \cap L^{\infty}(\Omega)$ and let $u_n \in H_0^1(\Omega)$ be the unique solution of

$$-\Delta u_n + V_n u_n + \lambda u_n = f$$

where $V_n = q_n^+ - q^- + iq_n^+$ and

$$q'_n = \begin{cases} n & \text{if} & q' > n \\ \\ q' & \text{if} & |q'| \leq n \\ \\ -n & \text{if} & q' \leq -n \end{cases}$$

The existence of u_n follows from a theorem of Lax-Milgram. Multiplying (16) by u_n we find

$$\|\mathbf{u}_{\mathbf{n}}\|_{\mathbf{H}^{1}} \leq \mathbf{c}$$

$$\int q_n^+ |u_n|^2 \le c .$$

Indeed let $u \in H_0^1(\Omega)$ with $u \ge 0$ a.e. on Ω ; let $u_n \in D(\Omega)$ be such that $u_n + u$ in $H^1(\Omega)$. We claim that $|u_n| + |u| = u$ in $H^1(\Omega)$ because $||u_n||_{H^1} = ||u_n||_{H^1}$ and $|u_n| \to |u|$ weakly in $H^1(\Omega)$. On the other hand $|u_n|$ can be smoothed by convolution and for fixed $n, \rho_{\epsilon} * |u_n| \to |u_n|$ in $H^1(\Omega)$ as $\epsilon \to 0$.

On the other hand we have

$$\Delta |u_n| \ge \text{Re}[\Delta u_n \text{ sign } \overline{u_n}] \text{ in } D'(\Omega)$$

which leads to

$$-\Delta |u_n| - q^-|u_n| + \lambda |u_n| \le |f|$$
 in $D'(\Omega)$.

Let $\psi \in H_0^1(\Omega)$ be the solution of

(19)
$$-\Delta \psi - q \psi + \lambda \psi = |f|.$$

It follows from Lemma 3.1 that

(20)
$$|u_n| \leq \psi$$
 a.e. on Ω .

By Theorem 2.3 we know that $\psi \in L^p(\Omega)$ for every $p \in [2,\infty)$. We extract a subsequence, denoted again by u_n such that $u_n + u$ weakly in $H^1_0(\Omega)$, $u_n + u$ a.e. on Ω . We see as in the proof of Theorem 2.1 (Step 1) that $(q_n^+ - q^-)u_n + qu$ in $L^1_{loc}(\Omega)$. Therefore we have only to verify that $q_n'u_n + q'u$ in $L^1_{loc}(\Omega)$. We distinguish two cases:

a)
$$q' \in L^{1+\epsilon}_{loc}(\Omega)$$
,

b)
$$q^{-\epsilon} L_{loc}^{(m/2)+\epsilon}(\Omega)$$
.

Case a) From (20) we deduce that $u_n \to u$ in every L^p space, $2 \le p < \infty$ and so $q'_n u_n \to q' u$ in $L^1_{loc}(\Omega)$.

Step 2: A + λ_1 is accretive. Let u \in D(A) and set T = Vu. We have T \in H⁻¹(Ω) \cap L¹_{loc}(Ω) and

Re T
$$\cdot \bar{\mathbf{u}} = \mathbf{q} |\mathbf{u}|^2 \ge -\mathbf{q}^- |\mathbf{u}|^2 \in L^1(\Omega)$$
.

It follows from Lemma 2.2 that $q|u|^2 \in L^1(\Omega)$ and

$$\int q|u|^2 = Re(T,u) = Re(Au + \Delta u,u).$$

Therefore

(21)
$$\operatorname{Re}(\lambda u, u) = \int |\operatorname{grad} u|^2 + \int q|u|^2 \ge -\lambda_1 \int |u|^2.$$

Step 3: D(A) is dense in $L^2(\Omega)$. Given $f \in L^2(\Omega) \cap L^{\infty}(\Omega)$ we solve for large number the equation

(22)
$$u_n + \frac{1}{n} A u_n = f$$
.

We shall prove that $u_n \to f$ in $L^2(\Omega)$ as $n \to \infty$ —and as a consequence D(A) is dense in $L^2(\Omega)$. By (21) we have

$$\int |u_n|^2 + \frac{1}{n} \int |gradu_n|^2 + \frac{1}{n} \int q|u_n|^2 = Re(f, u_n)$$
.

In particular we deduce that

(23)
$$\lim_{n\to\infty} \sup_{L^2} \left\| u_n \right\|_{L^2} \leq \left\| f \right\|_{L^2}$$

$$\frac{1}{n} \int q^+ |u_n|^2 \le C$$

(25)
$$\frac{1}{n} \int |\operatorname{gradu}_n|^2 \leq C.$$

Next we have (as in the proof of Step 1)

$$|u_n| - \frac{1}{n} \Delta |u_n| - \frac{1}{n} q^{-} |u_n| \le |f| \quad \text{in} \quad \mathcal{D}^{\bullet}(\Omega) \ .$$

On the other hand let $\psi \in H^1_0(\Omega)$ be the solution of

$$-\Delta \psi - q \psi + \lambda \psi = |f|$$

for some fixed $\lambda > \lambda_1$. Since $|u_n| \ge \lambda \left|\frac{u_n}{n}\right|$ for $n \ge \lambda$, we deduce from Lemma 3.1 that $\left|\frac{u_n}{n}\right| \le \psi$ a.e. Choose a subsequence, denoted again by u_n such that $u_n + u$ weakly in $L^2(\Omega)$, $\frac{1}{n}u_n + 0$ a.e. (this is possible since $\frac{1}{n}u_n + 0$ in $L^2(\Omega)$). For every $\varphi \in D(\Omega)$ we have

(26)
$$\int u_n \overline{\varphi} - \frac{1}{n} \int u_n \Delta \overline{\varphi} + \frac{1}{n} \int v u_n \overline{\varphi} = \int f \overline{\varphi} .$$

We claim that $\frac{1}{n}\int Vu_n \bar{\varphi} \to 0$ as $n \to \infty$. Indeed by (24) and (25) we have

$$\frac{1}{n} \mid \int q^+ u_n \bar{\phi} \mid \leq \frac{C}{\sqrt{n}} \quad \text{and} \quad \frac{1}{n} \mid \int q^- u_n \bar{\phi} \mid \leq \frac{C}{\sqrt{n}} \ .$$

Thus we have only to verify that $\frac{1}{n}\int q'u n^{\overline{\varphi}} \to 0$. We distinguish two cases:

- a) if $q \in L_{loc}^{(m/2)+\epsilon}(\Omega)$, we have $\psi \in L_{loc}^{\infty}(\Omega)$ and we deduce from the dominated convergence theorem that $\frac{1}{n} \int q' u_n \bar{\phi} \to 0$;
- b) if $q' \in L^{1+\varepsilon}_{loc}(\Omega)$ we use the fact that $\left|\frac{u}{n}\right| \leq \psi \in L^p(\Omega)$ for every $2 \leq p < \infty$ to deduce that $\frac{u}{n} + 0$ in $L^p(\Omega)$ and so $\frac{1}{n} \int q' u_n \tilde{\psi} + 0$.

In all the cases, we derive from (26) that

$$\int u\bar{\varphi} = \int f\bar{\varphi} \quad \forall \varphi \in D$$

and consequently u = f. We conclude using (23) that $u_n \to f$ in $\text{L}^2(\Omega)$.

Step 4: A is closable and $\bar{A} + \lambda_1$ is m-accretive. This is a standard fact, see e.g. Theorem 3.4 in [2].

Step 5: $u \in D(\overline{A})$ implies that $u \in H_0^1(\Omega)$, $q|u|^2 \in L^1(\Omega)$ and (14).

We already know (Step 2) that $v \in D(A)$ implies $q|v|^2 \in L^1(\Omega)$ and

(27)
$$\operatorname{Re}(Av,v) = \int |\operatorname{grad} v|^2 + \int q|v|^2.$$

Now let $u \in D(\overline{A})$ and let $u_n \in D(A)$ be such that $u_n \to u$, $Au_n \to \overline{A}u$. It follows from (27) applied to $v = u_n - u_m$ that $u_n \to u$ in $H_0^1(\Omega)$ and $\int q^+ |u_n - u|^2 \to 0$ (since u_n is a Cauchy sequence in $H_0^1(\Omega)$ and in $L^2(\Omega)$ with weight q^+). In particular $q|u|^2 \in L^1(\Omega)$ and (14) holds.

<u>Proof of Theorem 3.2</u>: Clearly $A \subseteq A_1$. Now let $u \in D(A_1)$ and let $\lambda > \lambda_1$. Set $f = A_1u + \lambda u$, and let u be the unique solution of

$$\bar{A}u^* + \lambda u^* = f$$
.

Thus, there exists a sequence $u_n^* \to u^*$ in $L^2(\mathbb{R}^m)$ with $u_n^* \in D(A)$ and

$$Au_n^* + \lambda u_n^* = f_n + f$$
 in $L^2(\mathbb{R}^m)$.

In particular we have

$$A_1(u_n^* - u) + \lambda(u_n^* - u) = f_n - f$$

and therefore

$$-\Delta \left| \mathbf{u}_{n}^{\bigstar} - \mathbf{u} \right| - \mathbf{q}^{-} \left| \mathbf{u}_{n}^{\bigstar} - \mathbf{u} \right| + \lambda \left| \mathbf{u}_{n}^{} - \mathbf{u} \right| \leq \left| \mathbf{f}_{n}^{} - \mathbf{f} \right| \quad \text{in} \quad \mathcal{D}^{\bullet}(\mathbf{R}^{m}) \ .$$

We deduce from Lemma 2.3 that $|u_n^* - u| \le \psi_n$ a.e. on \mathbb{R}^m where $\psi_n \in H^1(\mathbb{R}^m)$ is the solution of

$$-\Delta\psi_{\mathbf{n}} - \mathbf{q}^{\top}\psi_{\mathbf{n}} + \lambda\psi_{\mathbf{n}} = \left|\mathbf{f}_{\mathbf{n}} - \mathbf{f}\right| .$$

Hence $\|\psi_n\|_{\dot{H}^1} \to 0$ and in particular $u_n^* - u \to 0$ in $L^2(\mathbb{R}^m)$. It follows that $u^* = u$, that is $A_1 \subset \bar{A}$. We have $A \subset A_1 \subset \bar{A}$ and therefore A_1 is closable with $\overline{A_1} = \bar{A}$.

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SINGULAR COMPLEX POTENTIALS	6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(*) Haim Brezis and Tosio Kato	DAAG29-75-C-0024 MCS76-04655
Mathematics Research Center, University of 610 Walnut Street Madison, Wisconsin 53706	10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS Work Unit Number 1 - Applied Analysis
11. CONTROLLING OFFICE NAME AND ADDRESS See Item 18 below.	12. REPORT DATE August 1978 13. NUMBER OF PAGES
	17
14. MONITORING AGENCY NAME & ADDRESS(II different from Controlling Office)	15. SECURITY CLASS. (of this report) UNCLASSIFIED
	15a. DECLASSIFICATION/DOWNGRADING

Approved for public release; distribution unlimited.

17. DISTRIBUTION STATEMENT (of the abetract entered in Block 20, if different from Report)

18. SUPPLEMENTARY NOTES

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19. KEY WORDS (Continue on reverse side if necessary and identity by block number)

Schrodinger operator

Complex potentials

m-accretive operator

20. ABSTRACT (Continue on reverse elde if necessary and identify by block number)

Let $A = -\Delta + V(x)$ be a Schrödinger operator on an (arbitrary) open set

 $\Omega \subset \mathbb{R}^m$, where $V \in L^1_{loc}(\Omega)$ is a complex valued function. We consider the

"maximal" realization of A in $L^2(\Omega)$ under Dirichlet boundary condition,

that is

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20. ABSTRACT - cont'd.

$$D(A) = \{ u \in H_0^1(\Omega); Vu \in L_{loc}^1(\Omega) \text{ and } -\Delta u + Vu \in L^2(\Omega) \}.$$

When $\Omega = \mathbb{R}^m$ we also consider the operator

$$A_1 = -\Delta + V$$

with domain

$$D(A_1) = \{u \in L^2(\Omega); Vu \in L^1_{loc}(\Omega) \text{ and } -\Delta u + Vu \in L^2(\Omega)\}$$
.

A special case of our main results is:

Theorem: Let $m \geq 3$; assume that the function $\max\{-\text{Re V}, 0\}$ belongs to $L^{\infty}(\Omega) + L^{m/2}(\Omega)$ and also to $L^{(m/2)+\epsilon}_{\text{loc}}(\Omega)$ for some $\epsilon > 0$. Then A (resp. A_1) is closable and $\bar{A} + \lambda$ (resp. $\bar{A}_1 + \lambda$) is m-accretive for some real constant λ .